

#### 4.2 Brzina centra inercije krutog tela ( $V_i$ )

Vektor položaja centra inercije  $C_i$  krutog tela ( $V_i$ ) u odnosu na nepokretnu tačku O (vidi sl.4.1) određen je relacijom:

$$\overline{OC_i} = \vec{r}_i = \sum_{k=1}^i (\vec{\rho}_k + \xi_k \vec{e}_k q^k) + \vec{\tau}_i. \quad (4.22)$$

Prethodna relacija ukazuje na činjenicu da važi:

$$\vec{r}_i = \vec{r}_i(q^1, q^2, \dots, q^i), \quad (4.23)$$

odakle je

$$\vec{v}_{Ci} = \vec{v}_i = \sum_{\alpha=1}^i \vec{T}_{\alpha(i)} \dot{q}^\alpha \quad (4.24)$$

gde je  $\vec{T}_{\alpha(i)} = \partial \vec{r}_i / \partial q^\alpha$  pošto je prema (4.11)

$$\begin{aligned} \frac{\partial \vec{\rho}_k}{\partial q^\alpha} &= \vec{\xi}_\alpha \vec{e}_\alpha \times \vec{\rho}_k \quad \forall \alpha \leq k, \\ \frac{\partial \vec{\rho}_k}{\partial q^\alpha} &= 0 \quad \forall \alpha > k, \\ \frac{\partial \vec{e}_k}{\partial q^\alpha} &= \vec{\xi}_\alpha \vec{e}_\alpha \times \vec{e}_k \quad \forall \alpha \leq k, \\ \frac{\partial \vec{e}_k}{\partial q^\alpha} &= 0 \quad \forall \alpha > k. \end{aligned} \quad (4.25)$$

sledi (vidi (4.22)):

$$\begin{aligned} \vec{T}_{\alpha(i)} &= \vec{\xi}_\alpha \vec{e}_\alpha \times \left[ \sum_{k=\alpha}^i (\vec{\rho}_k + \xi_k \vec{e}_k q^k) + \vec{\tau}_i \right] + \xi_\alpha \vec{e}_\alpha \quad \forall \alpha \leq i, \\ \vec{T}_{\alpha(i)} &= 0 \quad \forall \alpha > i \end{aligned} \quad (4.26)$$

i

$$\vec{v}_i = \sum_{\alpha=1}^n \vec{T}_{\alpha(i)} \dot{q}^\alpha, \quad (4.27)$$

Poslednji izraz u matricnoj formi\* ima oblik:

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\* Gornji indeks ( $i$ ) označava da su koordinate odgovarajućeg vektora date u odnosu na lokalni koordinatni sistem  $C_i \xi_i \eta_i \zeta_i$  vezan za ( $V_i$ ). Specijalno, kada je  $i = 0$  u pitanju su koordinate date u odnosu na nepokretni koordinatni sistem  $Oxyz$ . Kada je u pitanju

$$\{\bar{v}_i^{(0)}\} = \sum_{\alpha=1}^n \{\bar{T}_{\alpha(i)}^{(0)}\} \dot{q}^\alpha, \quad (4.28)$$

gde je (vidi (4.25)):

$$\begin{aligned} \{\bar{T}_{\alpha(i)}^{(0)}\} &= \bar{\xi}_\alpha [A_{0,\alpha}] [e_\alpha^d] \left\{ \sum_{k=\alpha}^i [A_{\alpha,k}] (\{\bar{\rho}_k\} + \xi_k q^k \{\bar{e}_k\}) + [A_{\alpha,i}] \{\bar{\tau}_i\} \right\} + \\ &\quad + \bar{\xi}_\alpha [A_{0,\alpha}] \{\bar{e}_\alpha\} \quad \forall \alpha \leq i, \\ \{\bar{T}_{\alpha(i)}^{(0)}\} &= 0 \quad \forall \alpha > i \end{aligned} \quad (4.29)$$

Matrični zapis izraza (4.27) može da se dovede i na oblik:

$$\{\bar{v}_i^{(0)}\} = [E] \{\dot{q}\}, \quad (4.30)$$

gde je

$$[E] \in R^{3 \times n} \Rightarrow [E] = \left[ \{\bar{T}_{1(i)}^{(0)}\}; \{\bar{T}_{2(i)}^{(0)}\}; \dots; \{\bar{T}_{n(i)}^{(0)}\} \right], \quad (4.31)$$

$$\{\dot{q}\} \in R^{n \times 1} \Rightarrow \{\dot{q}\}^T = (\dot{q}^1, \dot{q}^2, \dots, \dot{q}^n). \quad (4.32)$$

#### 4.2.3 Brzina vrha hvataljke H robotskog sistema

**Primer 3.** Izvesti izraz za brzinu vrha hvataljke H robotskog sistema, slika 4.2.

Položaj vrha hvataljke je određen sa:

$$\bar{r}_H = \sum_{\alpha=1}^n (\bar{\rho}_{\alpha\alpha} + \bar{\xi}_\alpha \bar{e}_\alpha q^\alpha) \Rightarrow \bar{r}_H = \bar{r}_H(q^1, q^2, \dots, q^n), \quad (4.25)$$

tako da je, primenom iste procedure kao i u prethodnom primeru dobija

$$\bar{v}_H = \frac{d\bar{r}_H}{dt} = \sum_{\beta=1}^n \frac{\partial \bar{r}_H}{\partial q^\beta} \dot{q}^\beta = \sum_{\beta=1}^n \bar{\tau}_{\beta(n)} \dot{q}^\beta, \quad (4.26)$$

vektor vezan za telo ( $V_i$ ) (na primer  $\bar{\rho}_i = \overline{A_i B_i}$ ,  $A_i, B_i \in (V_i)$ ) tada u matričnom zapisu može da se izostavi gornji indeks ( $i$ ) ako se radi o koordinatama toga vektora, datim u odnosu na  $C_i \xi_i \eta_i \zeta_i$  (tj. o koordinatama datim u odnosu na lokalni koordinatni sistem za koje je vezan pomenuti vektor). Ova primedba se odnosi i na dualne objekte tih vektora.

gde je  $\bar{\tau}_{\beta(n)} = \frac{\partial \bar{r}_H}{\partial q^\beta}$  i on se razlikuje od kvazibaznog vektora  $\bar{T}_{\beta(n)}$  po odsustvu vektora  $\bar{\rho}_n$  i njegovih izvoda što se vidi iz sledećih izraza:

$$\bar{T}_{\beta(n)} = \bar{\xi}_\beta \bar{e}_\beta \times \left( \sum_{\alpha=\beta}^n (\bar{\rho}_{\alpha\alpha} + \bar{\xi}_\alpha \bar{e}_\alpha q^\alpha) + \bar{\rho}_n \right) + \bar{\xi}_\beta \bar{e}_\beta, \quad (4.27)$$

$$\bar{\tau}_{\beta(n)} = \bar{\xi}_\beta \bar{e}_\beta \times \left( \sum_{\alpha=\beta}^n (\bar{\rho}_{\alpha\alpha} + \bar{\xi}_\alpha \bar{e}_\alpha q^\alpha) \right) + \bar{\xi}_\beta \bar{e}_\beta. \quad (4.28)$$

### 4.3 Ubrzanje centra inercije krutog tela ( $V_i$ )

Ako se izraz (4.27) diferencira po vremenu dobijamo:

$$\bar{a}_{Ci} = \bar{a}_i = \sum_{\alpha=1}^i \bar{T}_{\alpha(i)} \ddot{q}^\alpha + \sum_{\alpha=1}^i \frac{d\bar{T}_{\alpha(i)}}{dt} \dot{q}^\alpha, \quad (4.33)$$

Kako je  $\bar{T}_{\alpha(i)} = \bar{T}_{\alpha(i)}(q^1, q^2, \dots, q^i)$  sledi:

$$\frac{d\bar{T}_{\alpha(i)}}{dt} = \sum_{\beta=1}^i \frac{\partial \bar{T}_{\alpha(i)}}{\partial q^\beta} \dot{q}^\beta. \quad (4.34)$$

Kako je  $\bar{T}_{\alpha(i)} = \partial \bar{r}_i / \partial q^\alpha$  (uz ispunjenje poznatih uslova o neprekidnosti i diferencijabilnosti funkcije  $\bar{r}_i = \bar{r}_i(t)$ ,  $i = 1, 2, \dots, n$ ) očigledno sledi:

$$\frac{\partial \bar{T}_{\alpha(i)}}{\partial q^\beta} = \frac{\partial^2 \bar{r}_i}{\partial q^\alpha \partial q^\beta} \Rightarrow \frac{\partial \bar{T}_{\alpha(i)}}{\partial q^\beta} = \frac{\partial \bar{T}_{\beta(i)}}{\partial q^\alpha}, \quad (4.35)$$

tako da je izraz (4.33) moguće dovesti na oblik:

$$\bar{a}_i = \sum_{\alpha=1}^i \bar{T}_{\alpha(i)} \ddot{q}^\alpha + \sum_{\alpha=1}^i \sum_{\beta=1}^i \bar{\Gamma}_{\alpha\beta(i)} \dot{q}^\alpha \dot{q}^\beta, \quad (4.36)$$

gde je

$$\bar{\Gamma}_{\alpha\beta(i)} = \frac{\partial \bar{T}_{\alpha(i)}}{\partial q^\beta}, \quad (4.37)$$

i

$$\bar{\Gamma}_{\alpha\beta(i)} = \bar{\Gamma}_{\beta\alpha(i)}, \quad (4.38)$$

U slučaju  $\alpha \leq \beta$  sledi (vidi (4.37) i (4.26)):

$$\bar{\Gamma}_{\alpha\beta(i)} = \bar{\xi}_\alpha \bar{e}_\alpha \times \frac{\partial}{\partial q^\beta} \left[ \sum_{k=\alpha}^i (\bar{\rho}_k + \xi_k \bar{e}_k q^k) + \bar{\tau}_i \right], \quad (4.39)$$

očigledno je da važi:

$$\frac{\partial}{\partial q^\beta} \left[ \sum_{k=\alpha}^i (\bar{\rho}_k + \xi_k \bar{e}_k q^k) + \bar{\tau}_i \right] = \frac{\partial}{\partial q^\beta} \left[ \sum_{k=1}^i (\bar{\rho}_k + \xi_k \bar{e}_k q^k) + \bar{\tau}_i \right], \quad (4.40)$$

ili (vidi (4.22) i imajući u vidu  $\bar{T}_{\alpha(i)} = \partial \bar{r}_i / \partial q^\alpha$ )

$$\frac{\partial}{\partial q^\beta} \left[ \sum_{k=\alpha}^i (\bar{\rho}_k + \xi_k \bar{e}_k q^k) + \bar{\tau}_i \right] = \bar{T}_{\beta(i)}. \quad (4.41)$$

Poslednji izraz daje relaciji (4.39) oblik:

$$\bar{\Gamma}_{\alpha\beta(i)} = \bar{\xi}_\alpha \bar{e}_\alpha \times \bar{T}_{\beta(i)} \quad \forall \alpha \leq \beta. \quad (4.42)$$

U slučaju  $\alpha > \beta$  sledi (vidi (4.37), (4.11))

$$\begin{aligned} \bar{\Gamma}_{\alpha\beta(i)} = & \bar{\xi}_\alpha \bar{\xi}_\beta (\bar{e}_\beta \times \bar{e}_\alpha) \times \left[ \sum_{k=\alpha}^i (\bar{\rho}_k + \xi_k \bar{e}_k q^k) + \bar{\tau}_i \right] + \\ & + \bar{\xi}_\alpha \bar{\xi}_\beta (\bar{e}_\beta \times \bar{e}_\alpha) + \bar{\xi}_\alpha \bar{e}_\alpha \times \left\{ \bar{\xi}_\beta \bar{e}_\beta \times \left[ \sum_{k=\alpha}^i (\bar{\rho}_k + \xi_k \bar{e}_k q^k) + \bar{\tau}_i \right] \right\}, \end{aligned} \quad (4.43)$$

što se nakon razvijanja dvostrukih vektorskih proizvoda može dovesti na oblik:

$$\bar{\Gamma}_{\alpha\beta(i)} = \bar{\xi}_\beta \bar{e}_\beta \times \left\{ \bar{\xi}_\alpha \bar{e}_\alpha \times \left[ \sum_{k=\alpha}^i (\bar{\rho}_k + \xi_k \bar{e}_k q^k) + \bar{\tau}_i \right] + \bar{\xi}_\alpha \bar{e}_\alpha \right\}, \quad (4.44)$$

ili, prema (4.26):

$$\bar{\Gamma}_{\alpha\beta(i)} = \bar{\xi}_\beta \bar{e}_\beta \times \bar{T}_{\alpha(i)} \quad \forall \alpha > \beta. \quad (4.45)$$

Prema (4.38) sledi da se pri sračunavanju vektora (4.37) može da koristi ili izraz (4.42) ili (4.45) a potom uzme u obzir činjenica da važi (4.38) (na primer, ako se određuje  $\bar{\Gamma}_{53(11)}$  može se iskoristiti izraz (4.42) prema kome se određuje  $\bar{\Gamma}_{53(11)}$  i potom uzme u obzir činjenica da važi:  $\bar{\Gamma}_{53(11)} = \bar{\Gamma}_{35(11)}$ ). Izraz (4.36) može da se napiše i u obliku:

$$\bar{a}_i = \sum_{\alpha=1}^n \bar{T}_{\alpha(i)} \ddot{q}^\alpha + \sum_{\alpha=1}^n \sum_{\beta=1}^n \bar{\Gamma}_{\alpha\beta(i)} \dot{q}^\alpha \dot{q}^\beta, \quad (4.46)$$

Pošto važi (vidi (4.26))

$$\bar{T}_{\alpha(i)} = 0 \quad \forall \alpha > i, \quad (4.47)$$

i\* (vidi (4.35)) i (4.37))):

$$\bar{\Gamma}_{\alpha\beta(i)} = 0 \quad \forall \sup(\alpha, \beta) > i, \quad (4.48)$$

Očigledno je da izrazi (4.42)) i (4.45) mogu da se napišu u jedinstvenoj formi

$$\bar{\Gamma}_{\alpha\beta(i)} = \bar{\xi}_a \bar{e}_a \times \bar{T}_{b(i)} \quad , \quad (4.49)$$

gde je \*\*

$$a = \inf(\alpha, \beta), \quad b = \sup(\alpha, \beta), \quad (4.50)$$

U matičnoj formi izraz (4.46) ima oblik:

$$\{\bar{a}_i^{(0)}\} = \sum_{\alpha=1}^n \{\bar{T}_{\alpha(i)}^{(0)}\} \ddot{q}^\alpha + \sum_{\alpha=1}^n \sum_{\beta=1}^n \{\bar{\Gamma}_{\alpha\beta(i)}^{(0)}\} \dot{q}^\alpha \dot{q}^\beta, \quad (4.51)$$

gde je (uzimamo da je  $\bar{\Gamma}_{\alpha\beta(i)}$  sračunat prema (4.42) a u slučaju  $\alpha > \beta$  iskorišćena je osobina simetrije po indeksima  $\alpha$  i  $\beta$ ):

$$\{\bar{\Gamma}_{\alpha\beta(i)}^{(0)}\} = \bar{\xi}_\alpha [A_{0,\alpha}] [e_\alpha^d] [A_{\alpha,\beta}] \{\bar{T}_{\beta(i)}^{(\beta)}\}, \quad (4.52)$$

gde je

$$\{\bar{T}_{\beta(i)}^{(\beta)}\} = \bar{\xi}_\beta [e_\beta^d] \left\{ \sum_{k=\alpha}^i [A_{\beta,k}] (\{\bar{p}_k\} + \xi_k q^k \{\bar{e}_k\}) + [A_{\beta,i}] \bar{\tau}_i \right\} + \xi_\beta \{\bar{e}_\beta\}, \quad (4.53)$$

#### 4.4 Ugaona brzina krutog tela ( $V_i$ )

Ugaona brzina tela ( $V_1$ ) koje je posredstvom cilindričnog (prizmatičnog) zgloba (1) (vidi sl.\*) vezano za postolje, ima ugaonu brzinu:

$$\bar{\omega}_1 = \bar{\xi}_1 \bar{e}_1 \dot{q}_1. \quad (4.54)$$

Kruto telo ( $V_2$ ) relativno se kreće u odnosu na telo ( $V_1$ ) (ili se obrće oko ose zgloba (2) vezane za telo ( $V_1$ ) ili se translatorno pravolinijski kreće duž ose zgloba (2) vezane za telo ( $V_1$ )). Relativna ugaona brzina tela ( $V_2$ ) u odnosu na ( $V_1$ ) iznosi:

$$\bar{\omega}_{2r} = \bar{\xi}_2 \bar{e}_2 \dot{q}^2, \quad (4.55)$$

Prema izrazu za apsolutnu brzinu tela ( $V_2$ ):

\*  $\sup(\alpha, \beta) = \alpha \quad \forall \alpha \geq \beta; \quad \sup(\alpha, \beta) = \beta \quad \forall \beta > \alpha$

\*\*  $\inf(\alpha, \beta) = \alpha \quad \forall \alpha \leq \beta; \quad \inf(\alpha, \beta) = \beta \quad \forall \alpha > \beta$

$$\vec{\omega}_2 = \vec{\omega}_{2p} + \vec{\omega}_{2r} . \quad (4.56)$$

gde je  $\vec{\omega}_{2p}$  -prenosna ugaona brzina tela ( $V_2$ ), uzimajući u obzir da je:

$$\vec{\omega}_{2p} = \vec{\omega}_1 , \quad (4.57)$$

sledi

$$\vec{\omega}_2 = \bar{\xi}_1 \bar{e}_1 \dot{q}^1 + \bar{\xi}_2 \bar{e}_2 \dot{q}^2 . \quad (4.58)$$

Polazeći od izraza za apsolutnu ugaonu brzinu tela ( $V_3$ ):

$$\vec{\omega}_3 = \vec{\omega}_{3p} + \vec{\omega}_{3r} , \quad (4.59)$$

Koristeći očigledne relacije:

$$\vec{\omega}_{3p} = \vec{\omega}_2 , \quad \vec{\omega}_{3r} = \bar{\xi}_3 \bar{e}_3 \dot{q}^3 , \quad (4.60)$$

dobija se:

$$\vec{\omega}_3 = \sum_{k=1}^3 \bar{\xi}_k \bar{e}_k \dot{q}^k . \quad (4.61)$$

Uopštavanjem prethodnih rezultata dolazi se do sledećeg izraza za ugaonu brzinu ( $\vec{\omega}_i$ ) krutog tela ( $V_i$ ):

$$\vec{\omega}_i = \sum_{k=1}^i \bar{\xi}_k \bar{e}_k \dot{q}^k \quad (4.62)$$

ili

$$\vec{\omega}_i = \sum_{k=1}^n \bar{\Omega}_{k(i)} \dot{q}^k , \quad (4.63)$$

u matricnoj formi relacija (4.62) glasi:

$$\{\vec{\omega}_i^{(0)}\} = \sum_{k=1}^i \bar{\xi}_k [A_{0,k}] \{\bar{e}_k\} \dot{q}^k , \quad (4.64)$$

ili

$$\{\vec{\omega}_i^{(0)}\} = [F] \{\dot{q}\} , \quad (4.65)$$

gde je

$$[F] \in R^{3 \times n} \Rightarrow [F] = \left[ \bar{\xi}_1 \{\bar{e}_1^{(0)}\} : \bar{\xi}_2 \{\bar{e}_2^{(0)}\} : \dots : \bar{\xi}_i \{\bar{e}_i^{(0)}\} \right] , \quad (4.66)$$

pri čemu je potrebno uzeti u obzir da važi

$$\{\vec{e}_k^{(0)}\} = [A_{0,k}] \{\vec{e}_k\}, \quad k = 1, 2, \dots, i. \quad (4.67)$$

#### 4.5 Ugaono ubrzanje krutog tela ( $V_i$ )

Diferenciranjem relacije (4.63) po vremenu dobija se izraz za ubrzanje tela ( $V_i$ ) u obliku:

$$\vec{\varepsilon}_i = \sum_{\beta=1}^n \vec{\Omega}_{\beta(i)} \ddot{q}^\beta + \sum_{\beta=1}^n \sum_{\alpha=1}^n \bar{\Lambda}_{\beta\alpha(i)} \dot{q}^\beta \dot{q}^\alpha, \quad (4.68)$$

gde je

$$\bar{\Lambda}_{\beta\alpha(i)} = \frac{\partial \vec{\Omega}_{\beta(i)}}{\partial \dot{q}^\alpha}, \quad (4.69)$$

odnosno sledi:

$$\bar{\Lambda}_{\beta\alpha(i)} = \vec{\xi}_\alpha \vec{e}_\alpha \times \vec{\Omega}_{\beta(i)} \Rightarrow \bar{\Lambda}_{\beta\alpha(i)} = \vec{\Omega}_{\alpha(\beta)} \times \vec{\Omega}_{\beta(i)}, \quad (4.70)$$

odakle zaključujemo da važi

$$\bar{\Lambda}_{\beta\alpha(i)} = 0 \quad \forall \alpha > \beta, \quad (4.71)$$

i, takodje,

$$\bar{\Lambda}_{\beta\alpha(i)} = 0 \quad \forall \beta > i. \quad (4.72)$$

Koristeći (4.70) i (4.71) izraz (4.68) možemo da dovedemo na oblik:

$$\vec{\varepsilon}_i = \sum_{\beta=1}^i \vec{\Omega}_{\beta(i)} \ddot{q}^\beta + \sum_{\beta=1}^i \sum_{\alpha=1}^{\beta} \vec{\Omega}_{\alpha(\beta)} \times \vec{\Omega}_{\beta(i)} \dot{q}^\alpha \dot{q}^\beta, \quad (4.73)$$

ili

$$\vec{\varepsilon}_i = \sum_{\beta=1}^i \vec{\Omega}_{\beta(i)} \ddot{q}^\beta + \sum_{\alpha=1}^i \sum_{\beta=\alpha}^i \vec{\Omega}_{\alpha(\beta)} \times \vec{\Omega}_{\beta(i)} \dot{q}^\alpha \dot{q}^\beta, \quad (4.74)$$

odakle sledi:

$$\vec{\varepsilon}_i = \sum_{\beta=1}^i \vec{\Omega}_{\beta(i)} \ddot{q}^\beta + \sum_{\alpha=1}^i \sum_{\beta=\alpha}^i \vec{\Omega}_{\alpha(\beta)} \times \vec{\Omega}_{\beta(i)} \dot{q}^\alpha \dot{q}^\beta, \quad (4.75)$$

ili

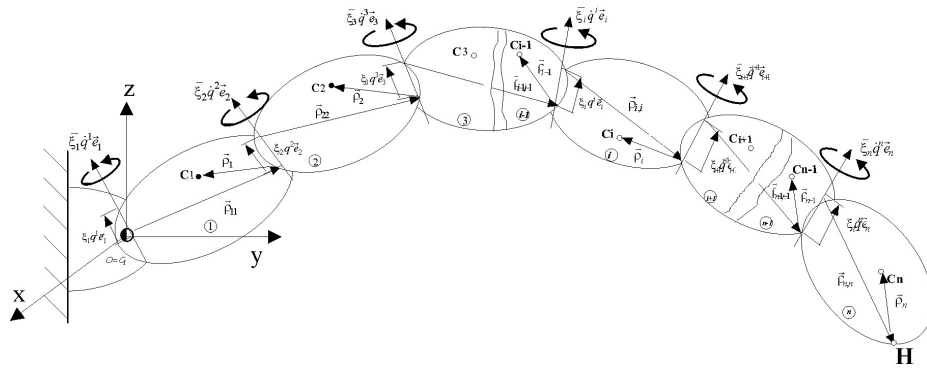
$$\vec{\varepsilon}_i = \sum_{\beta=1}^i \vec{\xi}_\beta \vec{e}_\beta \ddot{q}^\beta + \sum_{\alpha=1}^i \sum_{\beta=\alpha}^i \vec{\xi}_\alpha \vec{\xi}_\beta \vec{e}_\alpha \times \vec{e}_\beta \dot{q}^\alpha \dot{q}^\beta, \quad (4.76)$$

Poslednja relacija u matricnoj formi glasi:

$$\begin{aligned} \{\bar{\epsilon}_i^{(0)}\} &= \sum_{\beta=1}^i \bar{\xi}_{\beta} [A_{0,\beta}] \{\bar{e}_{\beta}\} \ddot{q}^{\beta} + \\ &+ \sum_{\alpha=1}^i \sum_{\beta=\alpha}^i \bar{\xi}_{\alpha} \bar{\xi}_{\beta} [A_{0,\alpha}] [e_{\alpha}^d] [A_{\alpha,\beta}] \{\bar{e}_{\beta}\} \dot{q}^{\alpha} \dot{q}^{\beta} \end{aligned} \quad (4.77)$$

**Primer za vezbanje.** Izvesti izraz za ugaonu brzinu centra inercije  $i$ -tog robotskog segmenta ( $V_i$ ) a zatim za robotski sistem sa pet stepeni slobode odrediti izraz za ugaonu brzinu poslednjeg –petog segmenta.

U odnosu na nepokretno postolje imamo apsolutno obrtanje samo u slučaju prvog segmenta, gde je  $q^1$  apsolutni ugao obrtanja. Prema tome sa  $\bar{\omega}_1$  je označena apsolutna ugaona brzina obrtanja prvog segmenta (4.1).



Slika 4.4a

$$\bar{\omega}_1 = \bar{\xi}_1 \bar{e}_1 \dot{q}^1 \quad (4.1)$$

Parametar  $\xi$  koji određuje da li je uočeni zglobov translatoran (T) ili rotacioni (R) je sada prema definiciji:

$$\xi_i = \begin{cases} 0, & \text{rotacioni zglobov} \\ 1, & \text{translatorni zglobov} \end{cases} \quad i = 1, 2, 3, \dots, n \quad (4.2)$$

odnosno:

$$\bar{\xi}_i = 1 - \xi_i, \quad i = 1, 2, 3, \dots, n \quad (4.3)$$



Telo ( $V_1$ ) vrši prenosno kretanje koje je ili translatorno ili rotaciono a telo ( $V_2$ ) relativno u odnosu na ( $V_1$ ) tj.:

$$\vec{\omega}_2 = \vec{\omega}_{2p} + \vec{\omega}_{2r}, \quad \vec{\omega}_{2p} = \vec{\omega}_1, \quad \vec{\omega}_{2r} = \bar{\xi}_2 \dot{q}^2 \vec{e}_2, \quad (4.4)$$

$$\vec{\omega}_2 = \bar{\xi}_1 \dot{q}^1 \vec{e}_1 + \bar{\xi}_2 \dot{q}^2 \vec{e}_2 \quad (4.5)$$

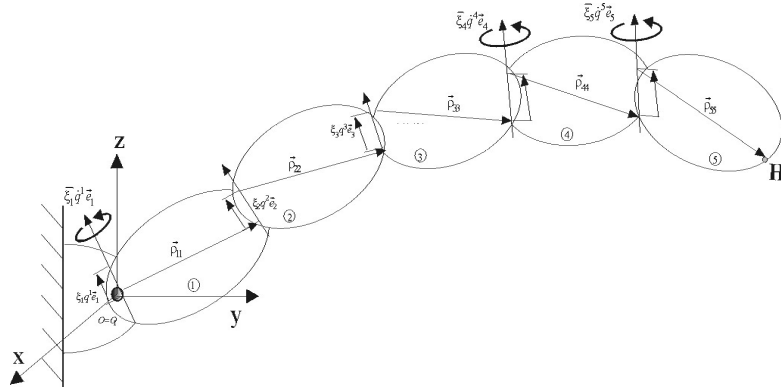
Na sličan način kao u slučaju drugog segmenta može se uočiti da i za treći segment važi:

$$\vec{\omega}_3 = \vec{\omega}_{3p} + \vec{\omega}_{3r}, \quad \vec{\omega}_{3p} = \vec{\omega}_2, \quad \vec{\omega}_{3r} = \bar{\xi}_3 \dot{q}^3 \vec{e}_3, \quad (4.6)$$

odnosno imajući u vidu relaciju (4.5) proizilazi da je:

$$\vec{\omega}_3 = \bar{\xi}_1 \dot{q}^1 \vec{e}_1 + \bar{\xi}_2 \dot{q}^2 \vec{e}_2 + \bar{\xi}_3 \dot{q}^3 \vec{e}_3 \quad (4.7)$$

$$\vec{\omega}_i = \sum_{\alpha=1}^i \bar{\xi}_{\alpha} \dot{q}^{\alpha} \vec{e}_{\alpha} \quad (4.8)$$



Slika 4.4b

U konkretnom slučaju za dati robotski sistem sa zglobovima RTTRR tj. imajući u vidu (4.2) i (4.3) dobija se  $\bar{\xi}_1 = 1, \bar{\xi}_2 = 0, \bar{\xi}_3 = 0, \bar{\xi}_4 = 1, \bar{\xi}_5 = 1$  odnosno na osnovu izraza (4.8) dobija se:

$$\vec{\omega}_5 = \dot{q}^1 \vec{e}_1 + \dot{q}^4 \vec{e}_4 + \dot{q}^5 \vec{e}_5 \quad (4.9)$$